Class 24: Disc dispersal ASTR 4008 / 8008, Semester 2, 2020

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Outline

- Demographics of late-stage discs
 - Accretion measures
 - Dust measures
 - Gas measures
- Models for disc dispersal
 - Viscous accretion
 - Photevaporation

Context: disc demographic surveys

- To understand how discs go away, want to be able to do big survey for presence of discs, and disc properties
- sources with imaging to check interpretation, resolve ambiguities

• Need to target a large range of stellar masses, and many star-forming regions of different age, to map out how disc properties change with mass and age

 Need for large samples means work mostly has to be done without imaging, because imaging is slow and expensive — at best can follow up some

• Three main strategies meet this requirements: (1) optical line diagnostics of accretion onto stellar photosphere, (2) IR / sub-mm continuum observations of dust continuum emission, (3) unresolved molecular line observations



Accretion measures The $H\alpha$ line

- Easiest disc signature to see in optical is stellar accretion
- Most common accretion signature is $H\alpha$ in emission
 - In most stars, $H\alpha$ seen in absorption signature of absorption by H atoms in n = 2 state in photosphere
 - Ha in emission requires gas above photosphere at T = 5 -10 kK at density high enough to get close to LTE
 - Material must have velocity spread of hundreds of km/s
- Model: emission from dense shock on stellar surface; can estimate accretion rate from brightness of line



Muzerolle+ 2005

Accretion demographics **Results from surveys**

- Can do large surveys of star-forming regions to find stars w/H α emission – these sources are relatively bright, so spectroscopy not too expensive
- Basic result: accretion fraction is ~50% at 3 Myr, ~10% at 10 Myr — e-folding time \approx 2 - 3 Myr
- Accretion rate scales with stellar mass roughly as dM/dt ~ M², but large scatter



Fedele+ 2010

IR / mm surveys Dust continuum emission

- IR / mm continuum traces dust in discs: NIR sensitive to dust near star (~1 AU), mm to dust far from star (~100 AU), intermediate wavelengths probe intermediate radii
- Basic result: e-folding time depends on wavelength / disc region
 - 2-3 Myr at NIR (~1 AU)
 - 4-6 Myr at MIR (~30 AU)
 - >10 Myr at mm (~100 AU)



Haisch+ 2001





Interpretation of dust continuum data What is the wavelength dependence telling us?

- Possibility 1: discs are cleared inside-out, so grains at ~1 AU hot enough to emit in NIR disappear earlier than those at ~30 AU which emit in MIR directly observed by imaging in some cases
- Possibility 2: longer wavelength emission at late times is not a signature of the original disc, but of a debris disc created when solid bodies created in first generation disc collide and produce dust
- Amount of dust required to produce longer wavelength emission is very small - 10^{-5} M $_{\oplus}$ sufficient (about mass of Vesta) – so possibility 2 plausible
- May be a combination of both



Dust vs accretion data Very confusing

- Accretion and hot dust decrease in time with similar e-folding
- However, little object-by-object correlation between dust holes or dust mass and accretion rates
- Plenty of objects with > 10 AU dust cavities still have gas accretion rates > 10^{-8} M $_{\odot}$ / yr



Ercolano & Pascucci 2017



Gas measures Challenges and status

- statistical studies
- Unresolved observations cheaper, and can use spectra to make inferences about presence of gas close to star — revealed by high Doppler velocities
- However, optical depth effects make results difficult to interpret for abundance species like CO, luminosity not just proportional to mass
- Gas clearing times seem roughly consistent with dust clearing times, but hard to be more specific given poor statistics and modelling uncertainties

Resolved gas observations very expensive, so sample size still small - no



Viscous accretion **Dispersal mechanism I**

- Disc could disappear just by all accreting onto star at a few Myr
- Mass budget reasonable: observed accretion rates ~10⁻⁸ 10⁻⁹ M_{\odot} / yr can drain a 0.01 M_{\odot} disc in ~1 - 10 Myr, about observed lifetime
- However, disc evolution obeys $\frac{\partial \Sigma}{\partial t} = \frac{3}{\varpi} \frac{\partial}{\partial \varpi} \left[\varpi^{1/2} \frac{\partial}{\partial \varpi} \left(\nu \Sigma \varpi^{1/2} \right) \right]$
- Generic solution for power law viscosity $\nu = \nu_1 (\varpi / \varpi_1)^{\gamma}$ is:
- power law in time \rightarrow disc never develops an inner hole

 $\Sigma = \frac{C}{3\pi\nu_{1}x^{\gamma}}T^{-(5-2\gamma)/(4-2\gamma)} \exp\left(-\frac{x^{2-\gamma}}{T}\right), \qquad x = \frac{\varpi}{\pi\tau_{1}}, \quad T = \frac{3(2-\gamma)\nu_{1}}{\pi\tau_{1}^{2}}t + 1$

• Problem: for small x (inner disc), surface density decreases uniformly as a

Viscous accretion Advantages and problems

- does not easily create holes, as observations seem to require
- Need some mechanism to produce sharp features in dust or gas distribution
- However, accretion clearly does occur, as shown by the H α data, and observed accretion rates are sufficient to drain fair amount of mass from disc
- Most promising avenue is accretion combined with some other mechanism to produce sharp features

• Problem with viscous accretion model is generic: anything based on viscosity



Dispersal mechanisms **II. Photoevaporation**



Ercolano & Pascucci 2017



Photoevaporation The basics

- Young stars produce FUV (11 13.6 eV) and EUV (13.6 eV ~100 eV) thermally, and soft X-rays (~100 eV - few keV) via non-thermal processes in the magnetosphere
- The radiation heats a surface layer of the disc to ~few hundred K (FUV), ~10⁴ K (EUV), or hotter (X-ray)
- If the sound speed in the heated gas exceeds the escape speed from the location of the hot gas, thermal pressure will launch a wind off the disc
- This mechanism naturally produces a sharp feature in the disc: mass is only lost at sufficiently large radii

Photoevaporation **Example calculation for EUV**

- probably the least important of the three mechanisms
- Basic idea: wind flows where escape speed ≤ 10 km/s: $\varpi \gtrsim GM/c_s^2 \sim 9(M/M_{\odot})$ au
- base of wind obey $\Phi \sim \alpha_B n_0^2 \varpi^3 =$
- Wind flows out at ~c_s, so

 $\dot{M} \sim n_0 \mu m_{\rm H}$

EUV is simplest because chemistry and temperature are simple, though it is

• If ionising luminosity of star is Φ , ionisation balance requires that density n₀ at

$$\Rightarrow n_0 \sim \sqrt{\frac{\alpha_B}{\alpha_B}} \frac{c_s}{(GM)^3}$$
$$c_s \sim \sqrt{\frac{\Phi}{\alpha_B}} \frac{c_s^4 \mu m_{\rm H}}{(GM)^{3/2}}$$

• Observed Φ values give ~10⁻¹⁰ M_{\odot} / yr for EUV, but FUV and X-rays stronger

Possible scenario Still not well understood

- stop gas
- the gap preventing material at larger radii from getting in
- inner disc, evaporation rate rises, disc is removed fairly quickly

Photoeavporation produces a gap in discs, perhaps aided by massive planets

 The gaps act as pressure traps that prevent dust grains from moving through; depending on evaporation rate and planet properties, may or may not also

• Material interior to the gap accretes onto star, and is not replenished due to

Once the inner rim of the gap is exposed to direct starlight by accretion of the