- 1. It is sometimes convenient to write down a simple analytic expression for the rate of cooling in the neutral ISM.
  - (a) We will start with C<sup>+</sup> cooling. Make the following assumptions: (1) only collisions between C<sup>+</sup> and H are significant; (2) the collisional *de*-excitation rate coefficient for such collisions,  $k_{C^+,de}$  is a constant that is independent of temperature; (3) the density is much less than the critical density of C<sup>+</sup>; (4) all carbon is C<sup>+</sup>. Under these assumptions, give a simple analytic formula for the C<sup>+</sup> cooling rate in gas with H number density  $n_{\rm H}$  and C/H ratio  $\delta_{\rm C}$ .
  - (b) Plot the cooling function you have just derived as a function of temperature. Use  $\delta_{\rm C} = 10^{-4}$ , and you can get a numerical value for  $k_{\rm C^+,de}$  from the Leiden Atomic and Molecular Database.
  - (c) Produce a similar analytic expression and plot for cooling by O, but for O assume that the de-excitation rate coefficients for the two upper states scale with temperature as  $k_{\text{O,de},0}\sqrt{T/T_0}$ .
- 2. One additional feature we can add to our model of vertical hydrostatic equilibrium is a layer of dense material near the midplane, which might for example represent dense molecular clouds. Suppose that we have a layer of atomic hydrogen with total surface density  $\Sigma_g$  and constant velocity dispersion  $\sigma_g$ , and also an infinitesimally thin layer of material with surface density  $\Sigma_{\rm mp}$  located at the midplane.
  - (a) Write down the equation of hydrostatic balance including the midplane material.
  - (b) Make the same transformation from  $\rho_q$  to  $s_q$  as your variable that we did in class.
  - (c) Solve in the limits  $\Sigma_{\rm mp} \ll \Sigma_g$  and  $\Sigma_{\rm mp} \gg g$ .
  - (d) Show that, for  $z \to \infty$ , the solution is the same in both limits, as long as  $\Sigma_{\rm g} + \Sigma_{\rm mp}$  remains constant.