Srid-Lused hydrodynamics

• Time derivative

Forward Fuller method (simplest):  $\partial_t q \longrightarrow \frac{q^{m+1} - q^m}{n+1}$  with  $\Delta t = t^{n+1} - t^n$ • Think of single ODE :  $\dot{q} = \partial_t q = f(q, t)$  $\implies \frac{q^{n+1}-q^n}{\lambda \neq} = f(q^n, \xi^n)$ (explicit method) (1st-order method) · becomes instable if at is chosen too big · could be made higher-order (Ringe-Kutta, Adams) · Spahial derivative  $\partial_{x} q = \lim_{\Delta x \to 0} \frac{q(x + \Delta x) - q(x)}{\Delta x}$ · in discretised form : i - 1 i i + 1 $\partial_x q \rightarrow \frac{q_{i+1} - q_{i-1}}{2Ax}$ (2nd order accurate) (Centred Difference)

Advection

- Start with simpler form of hydro eq., where u = const.=>  $\partial_t q + u \partial_x q = 0$
- · Insert time- und space-derivatives :

Justead upwind scheme:  

$$q_{i}^{n+1} = q_{i}^{n} - \frac{\Delta t}{\Delta x} u \left(q_{i-1}^{n} - q_{i}^{n}\right) (u > 0)$$
  
(Upwind Difference Scheme) = stable

• uprind schene stable, but diffusive 
$$\longrightarrow$$
  
numerical  
diffusion.  
• Diffusion eq. :  $\frac{\partial}{\partial t}q - D = \frac{\partial^2}{\partial x^2}q = D$ 

discretised version:



$$\frac{q_{i}-q_{i-1}}{\Delta x} = \frac{q_{i+1}-q_{i-1}}{2\Delta x} - \Delta x \frac{q_{i+1}-2q_{i}+q_{i-1}}{2\Delta x^{2}}$$
centred diff.  
= Uprind scheme is  
the centred difference  
Scheme with diffusion.  

$$\frac{q_{i+1}-2q_{i}+q_{i-1}}{2\Delta x^{2}}$$
with a diffusion  

$$\frac{q_{i+1}-2q_{i}+q_{i-1}}{2\Delta x^{2}}$$
diffusion  

$$\frac{q_{i+1}-2q_{i}+q_{i-1}}{2\Delta x^{2}}$$