

# ASTR4012/ASTR8002

## Astrophysical Gas Dynamics

### Assignment 4 – exam assignment

Christoph Federrath

due Tuesday, October 13, 2020  
(no extensions possible – exam assignment)

## 1 Validity of the gas/fluid approximation

In the following, calculate and discuss the conditions for approximating a collection of particles as a gas or fluid. Order-of-magnitude calculations are sufficient.

- (a) Calculate the particle mean free path ( $\lambda$ ) of air on Earth (typical density of air near the surface, typical temperature  $\sim 300$  K, etc.). Write down the size scales over which air can be treated as a gas and over which it cannot.  
(2 points)
- (b) What is particle mean free path in the solar corona (density  $\rho \sim 10^{-16}$  g cm $^{-3}$ )? Start from this density and make use of the mean particle weight in the solar corona. Briefly discuss (2–3 sentences) what the main difference is in computing  $\lambda$  for air and  $\lambda$  in the solar corona.  
(2 points)
- (c) Calculate the mean free path in a galaxy cluster, i.e., in the intra-cluster medium (ICM). Compare this mean free path to the typical size of a galaxy cluster and discuss briefly (again 2–3 sentences) the validity of the fluid/gas approximation in the ICM.  
(2 points)
- (d) What is the particle mean free path in a molecular cloud? Compare two cases: (1) assuming collisions of molecular hydrogen molecules if there are no charges and/or no magnetic fields present, and (2) for the case where the cross section for collisions would be primarily determined by electrostatic interactions.  
(1 point)
- (e) What would you do if you want to describe the dynamics of a system in which particle collisions are extremely rare events? Only write 1–2 sentences.  
(1 point)

## 2 Rankine-Hugoniot shock jump conditions

- (a) Find the compression ratio  $r \equiv \rho_2/\rho_1$  at a plane shock front in terms of the upstream Mach number  $\mathcal{M} = v_1/c_1$ , where  $c_1 = (\gamma P_1/\rho_1)^{1/2}$  is the upstream sound speed (subscripts 1 and 2 refer to the pre-shock and post-shock quantities, respectively). Assume that the gas flows along the shock normal.

*Hint: Use the hydrodynamic equations in conservation form to identify the flux of mass, momentum and energy entering and leaving the shock front.*

(5 points)

- (b) A supernova remnant shock moves at  $\sim 5,000$  km/s into a part of the interstellar medium with temperature of  $\sim 10^4$  K. Estimate the downstream (post-shock) temperature, assuming the medium consists of ionised hydrogen in full thermodynamic equilibrium.

(3 points)

- (c) What would be the approximate temperature behind a shock front moving at 50,000 km/s, if full thermodynamic equilibrium were established? Why does this approximation fail?

(2 points)

### 3 Evolution of a supernova remnant

Assume first the supernova remnant in its adiabatic phase: all the mass of the remnant is concentrated in a thin shell located at the position of the shock at radius  $r = r_s(t)$ , where  $M \approx (4\pi/3)r^3\rho_1$ , with  $\rho_1$  being the density of the ambient (interstellar) medium. Furthermore, the pressure  $P(t)$  interior to the shock can be considered uniform and the equation of motion for the thin shell is given by

$$\frac{d}{dt}(M\dot{r}_s) = 4\pi r_s^2 P, \quad (1)$$

where  $\dot{r}_s = dr_s/dt$ .

- (a) Use the jump conditions for a strong shock of an ideal gas with an adiabatic index  $\gamma = 5/3$  to estimate the thickness  $\Delta r$  of the shell in terms of  $r_s$  (assume the shell density equals the post-shock density).

(2 points)

- (b) Given that in the Sedov phase, the total internal energy of the gas in the remnant equals 80% of the explosion energy  $E_{\text{SN}}$ , show that the equations of motion have a solution of the form

$$r_s = At^\alpha. \quad (2)$$

Find the constants  $A$  and  $\alpha$ .

(4 points)

- (c) Assume now that the shell cools rapidly. Because the cooling rate of the gas is proportional to the density squared, there is a phase in the evolution when the thermal energy of the freshly shocked gas can no longer be shared evenly throughout the remnant, as it was the case in the energy-driven phase discussed above. Most of the energy is radiated away and a cool dense shell forms around the still hot interior. To a good approximation, the interior can be described as a hot adiabatic gas bubble of constant mass with an equation of state as before,  $P \propto \rho^\gamma$  with  $\gamma = 5/3$ . The evolution of the blast wave is now driven by the adiabatic expansion of the bubble. Show that this pressure-driven ‘snow plow phase’ admits again a solution of the form

$$r_s \propto t^\beta, \quad (3)$$

and find the index  $\beta$ .

*Hint: the equation of motion of the shell, Eq. (1), still applies in this phase.*

(4 points)

## 4 Magnetic breaking of a rotating gas disc

Consider a uniform gaseous disc of density  $\rho_{\text{cl}}$  and half-thickness  $Z$  rotating rigidly with an initial angular velocity  $\Omega_0$ . Furthermore assume the disc is threaded by a magnetic field  $\mathbf{B}$  of strength  $B_0$ , initially uniform and parallel to the rotation axis of the disc. The magnetic field couples the disc ( $|z| \leq Z$ ) with the external medium ( $|z| > Z$ ) of density  $\rho_{\text{ext}}$ , which is initially at rest. Use cylindrical coordinates  $(R, \phi, z)$ , assume symmetry around the rotation axis, i.e.,  $\partial\phi = 0$ , and assume ideal MHD.

- (a) For this system, derive the evolution equation for the angular velocity  $\Omega$ , where  $v_\phi = R\Omega$ , and the evolution equation for the toroidal component of the magnetic field,  $B_\phi$ . In your derivations, assume that  $v_R \ll v_\phi$ ,  $v_z \ll v_\phi$ , and  $B_R \ll B_0$ , allowing you to drop 2nd-order terms in the Euler and induction equation. The final results should be

$$\frac{\partial\Omega}{\partial t} = \frac{B_0}{4\pi R\rho} \frac{\partial B_\phi}{\partial z}, \quad \text{and} \quad (4)$$

$$\frac{\partial B_\phi}{\partial t} = B_0 R \frac{\partial\Omega}{\partial z}. \quad (5)$$

(4 points)

- (b) Show that the evolution of the external medium can be expressed by the wave equation

$$\frac{\partial^2\Omega}{\partial t^2} = v_{\text{A,ext}}^2 \frac{\partial^2\Omega}{\partial z^2}, \quad (6)$$

with the Alfvén speed  $v_{\text{A,ext}} = B_0/(4\pi\rho_{\text{ext}})^{1/2}$  in the external medium.

(2 points)

- (c) Derive the evolution equation of the angular velocity at the surface of the disc ( $|z| = Z$ ), using the torque per unit area  $N = RB_0B_\phi/(4\pi)$  that the magnetic field exerts on the surface of the disc. The result should be

$$\frac{\partial^2\Omega_{\text{cl}}}{\partial t^2} = \frac{1}{Z} \frac{\rho_{\text{ext}}}{\rho_{\text{cl}}} v_{\text{A,ext}}^2 \frac{\partial\Omega}{\partial z} \Big|_{|z|=Z} \quad (7)$$

*Hint: The full torque is  $I\partial\Omega/\partial t$  with the moment of inertia  $I$  for a cylinder.*

(3 points)

- (d) Combine Equations (6) and (7) to calculate the spin-down time of the disc.

*Hint: Use the solution of Equation (6) at the disc surface and make the ansatz  $\Omega \propto \exp(-\alpha t + \beta z)$  to determine the constants  $\alpha$  and  $\beta$ .*

(3 points)

**Please submit your solutions via Turnitin by the assignment deadline. Please note that this is the exam assignment, so absolutely no extensions are possible!**