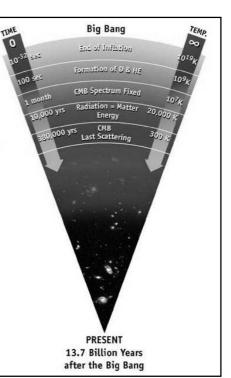
WMAP CMB Conclusions

- A flat universe with a scale-invariant spectrum of adiabatic Gaussian fluctuations, with re-ionization, is an acceptable fit to the WMAP data.
- The correlations of polarisation and the acoustic peaks imply the initial fluctuations were primarily adiabatic (the primordial ratios of dark matter/photons & baryons/photons do not vary spatially).
- The initial fluctuations are consistent with a Gaussian field, as expected from most inflationary models.

WMAP's cosmic timeline

- CMB last scattering surface: t_{dec} = 379 \pm 8 kyr (z_{dec} = 1089 \pm 1)
- Epoch of re-ionization:
 t_r = 100 400 Myr (95% c.l.)



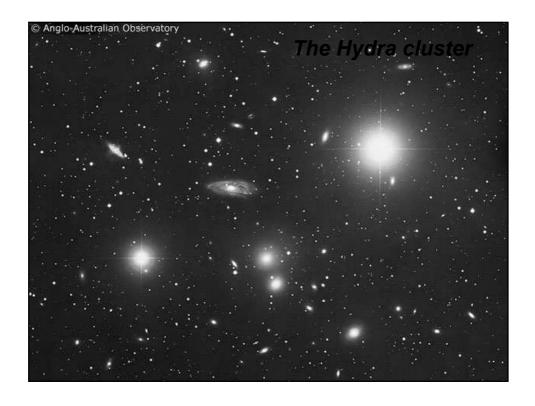
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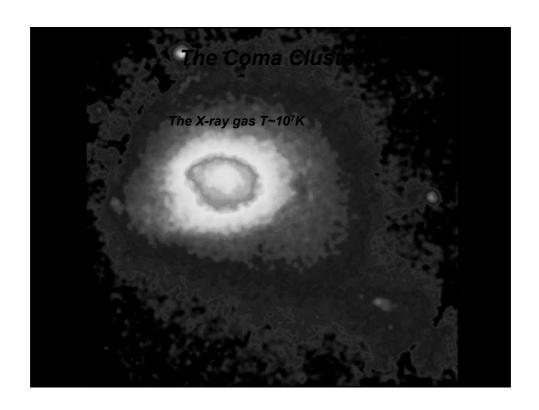
Large Scale Structure

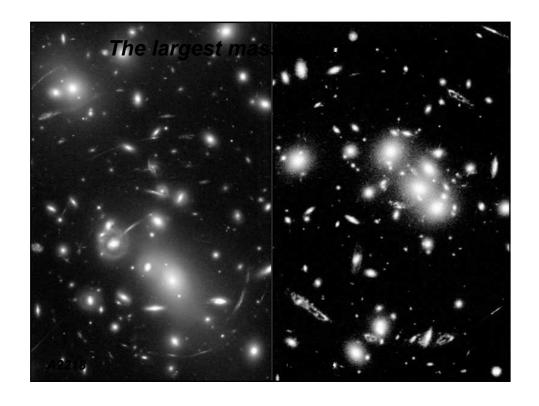
With Thanks to Matthew Colless, leader of Australian side of 2dF redshift survey.



The Local Group

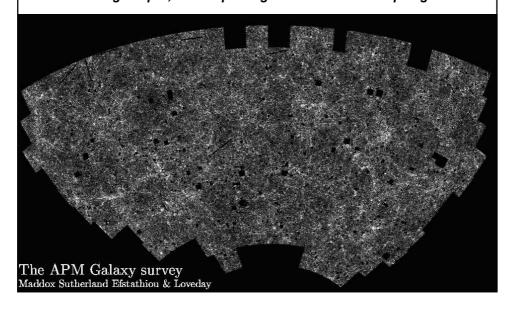






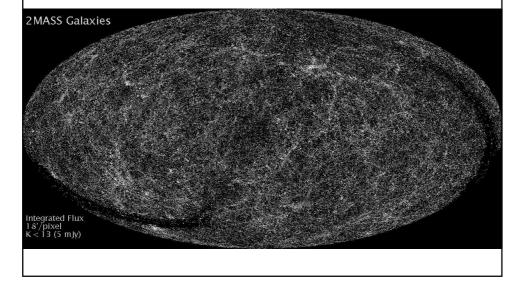
Large-scale structure in the local universe

Going deeper, 2x10⁶ optical galaxies over 4000 sq. deg.



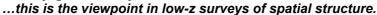
Large-scale structure in the local universe

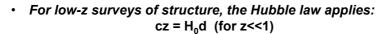
The 10⁶ near-infrared brightest galaxies on the sky



Redshift surveys

- A z-survey is a systematic mapping of a volume of space by measuring redshifts: $z = \lambda_1/\lambda_0 - 1 = a-1$
- Redshifts as distance coordinates...
 H₀D₁ = c(z+(1-q₀)z²/2+...)

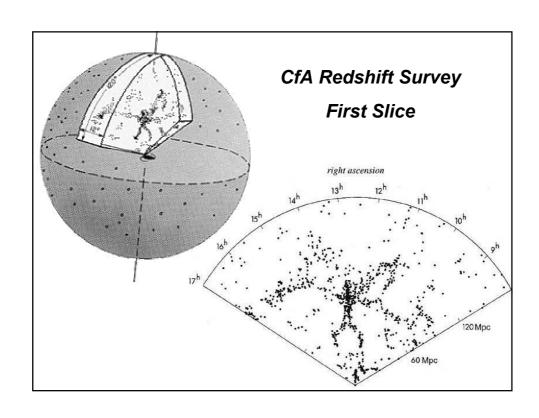


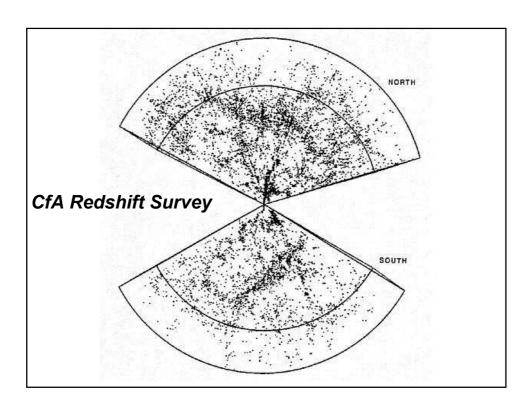


- For pure Hubble flow, redshift distance = true distance,
 i.e. s=r, where s and r are conveniently measured in km/s.
- But galaxies also have 'peculiar motions' due to the gravitational attraction of the surrounding mass field, so the full relation between z-space and real-space coordinates is : $s = r + v_p \cdot r/r = r + v_p$ (for s<<c)

Uses of z-surveys

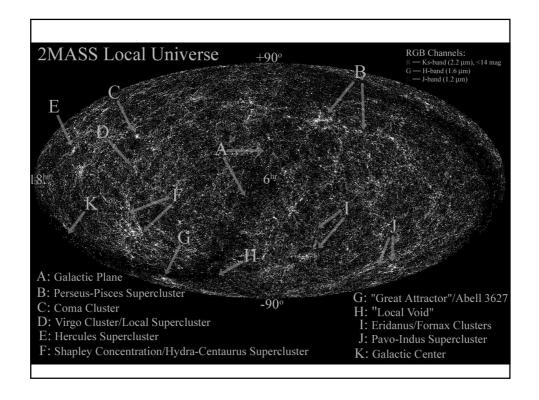
- Three (partial) views of redshift:
 - z measures the distance needed to map 3D positions
 - z measures the look-back time needed to map histories
 - cz-H₀d measures the peculiar velocity needed to map mass
- · Three main uses of z-surveys:
 - to map the large-scale structures, in order to...
 - · do cosmography and chart the structures in the universe
 - · test growth of structure through gravitational instability
 - · determine the nature and density of the dark matter
 - to map the large-velocity field, in order to
 - · `see' the mass field through its gravitational effects
 - to probe the history of galaxy formation, in order to...
 - characterise the galaxy population at each epoch
 - determine the physical mechanisms by which the population evolves





Cosmography

- The main features of the local galaxy distribution include:
 - Local Group: Milky Way, Andromeda and retinue.
 - Virgo cluster: nearest significant galaxy cluster, LG→Virgo.
 - Local Supercluster (LSC): flattened distribution of galaxies cz<3000 km/s; defines supergalactic plane (X,Y,Z).
 - 'Great Attractor': LG/Virgo→GA, lies at one end of the LSC, (X,Y,Z)=(-3400,1500,±2000).
 - Perseus-Pisces supercluster: (X,Y,Z)=(+4500, ±2000,-2000), lies at the other end of the LSC.
 - Coma cluster: nearest very rich cluster, (X,Y,Z)=(0,+7000,0);
 a major node in the 'Great Wall' filament.
 - Shapley supercluster: most massive supercluster within z<0.1, at a distance of 14,000 km/s behind the GA.
 - Voids: the Local Void, Sculptor Void, and others lie between these mass concentrations.
- Yet larger structures are seen at lower contrast to >100 h⁻¹ Mpc.



Evolution of Structure

• The goal is to derive the evolution of the mass density field, represented by the dimensionless density perturbation:

$$\delta(x) = \rho(x)/<\rho> - 1$$

- The framework is the growth of structures from 'initial' density fluctuations by gravitational instability.
- Up to the decoupling of matter and radiation, the evolution of the density perturbations is complex and depends on the interactions of the matter and radiation fields - 'CMB physics'
- After decoupling, the linear growth of fluctuations is simple and depends only on the cosmology and the fluctuations in the density at the surface of last scattering - 'large-scale structure in the linear regime'.
- As the density perturbations grow the evolution becomes nonlinear and complex structures like galaxies and clusters form -'non-linear structure formation'. In this regime additional complications emerge, like gas dynamics and star formation.

The power spectrum

• It is helpful to express the density distribution $\delta(\mathbf{r})$ in the Fourier domain:

$$\delta(k) = V^{-1} \int \delta(r) e^{ik \cdot r} d^3r$$

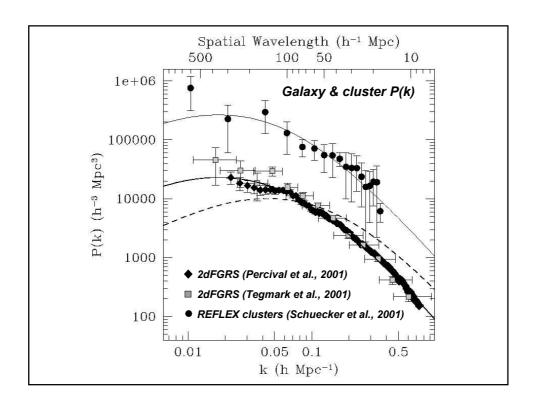
 The power spectrum (PS) is the mean squared amplitude of each Fourier mode:

$$P(k) = < |\delta(k)|^2 >$$

- Note P(k) not P(k) because of the (assumed) isotropy of the distribution (i.e. scales matter but directions don't).
- P(k) gives the power in fluctuations with a scale $r=2\pi/k$, so that k = (1.0, 0.1, 0.01) Mpc⁻¹ correspond to r≈ (6, 60, 600) Mpc.
- The PS can be written in dimensionless form as the variance per unit ln k:

$$\Delta^{2}(k) = d < |\delta(k)|^{2} > /dlnk = (V/2\pi)^{3} 4\pi k^{3} P(k)$$

- e.g. $\Delta^2(k)$ =1 means the modes in the logarithmic bin around wavenumber k have rms density fluctuations of order unity.



The correlation function

• The autocorrelation function of the density field (often just called 'the correlation function', CF) is:

$$\xi(r) = \langle \delta(x)\delta(x+r) \rangle$$

• The CF and the PS are a Fourier transform pair:

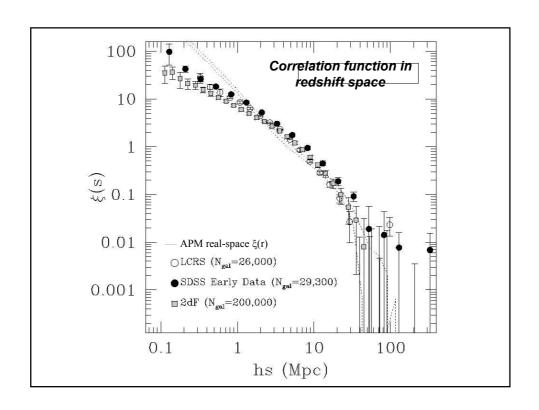
$$\xi(r) = V/(2\pi)^{3} \int |\delta_{k}|^{2} \exp(-ik \cdot r) d^{3}k$$

= $(2\pi^{2})^{-1} \int P(k)[(\sin kr)/kr] k^{2} dk$

- Because P(k) and $\xi(r)$ are a Fourier pair, they contain precisely the same information about the density field.
- When applied to galaxies rather than the density field, ξ(r) is often referred to as the 'two-point correlation function', as it gives the excess probability (over the mean) of finding two galaxies in volumes dV separated by r:

$$dP = \rho_0^2 [1 + \xi(r)] d^2V$$

- By isotropy, only separation r matters, and not the vector r.
- Can thus think of $\xi(r)$ as the mean over-density of galaxies at distance r from a random galaxy.



Gaussian fields

- A Gaussian density field has the property that the joint probability distribution of the density at any number of points is a multivariate Gaussian.
- Superposing many Fourier density modes with <u>random</u> <u>phases</u> results, by the central limit theorem, in a Gaussian density field.
- A Gaussian field is fully characterized by its mean and variance (as a function of scale).
- Hence and P(k) provide a complete <u>statistical</u> description of the density field if it is Gaussian.
- Most simple inflationary cosmological models predict that Fourier density modes with different wavenumbers are independent (i.e. have random phases), and hence that the initial density field will be Gaussian.
- Linear amplification of a Gaussian field leaves it Gaussian, so the large-scale galaxy distribution should be Gaussian.

The initial power spectrum

- Unless some physical process imposes a scale, the initial PS should be scale-free, i.e. a power-law, $P(k) \propto k^n$.
- The index n determines the balance between large- and small-scale power, with rms fluctuations on a mass scale M given by: $\delta_{rms} \propto M^{\cdot (n+3)/6}$
- The 'natural' initial power spectrum is the power-law with n=1 (called the Zel'dovich, or Harrison-Zel'dovich, spectrum).
- The $P(k) \propto k^1$ spectrum is referred to as the scale-invariant spectrum, since it gives variations in the gravitational potential that are the same on all scales.
- Since potential governs the curvature, this means that spacetime has the same amount of curvature variation on all scales (i.e. the metric is a fractal).
- In fact, inflationary models predict that the initial PS of the density fluctuations will be <u>approximately</u> scale-invariant.